# BeppoSAX observations of the quasar Markarian 205

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**Abstract.** We present the first BeppoSAX observation (0.1 to 220 keV) of the quasar Mrk 205. We have searched for the unusual Fe line profile claimed in the XMM-Newton spectrum which has been widely discussed in recent literature. We find no evidence for a broad, ionized Fe line component in our data. We detect for the first time a Compton hump in this object. Besides, when this component is included in the fit, the line strength diminishes, in agreement with a recent re-analysis of the XMM-Newton data, but with better constraints on the reflection component thanks to the PDS instrument (15-220 keV). We interpret this fact as another indication for illumination of a distant and cold material rather than reprocessing in the highly ionized inner parts of an accretion disk.

We cannot constrain the presence of a high energy cutoff but we confirm the existence of a variable soft excess (one year timescale).

Key words. Galaxies: active – quasars: individual: Mrk 205 – X-rays: galaxies

#### 1. Introduction

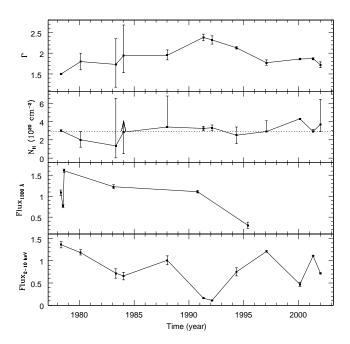
Discovered by Weedman (1970), the quasar Mrk 205 is seen through the southern spiral arm of the nearby spiral, barred galaxy NGC 4319 at a projected distance of roughly 40'' from its center. From its radio luminosity at 6 cm, Mrk 205 is classified as a radio-quiet quasar (Rush et al. 1996a), but it is still not clear whether it belongs to the radio-quiet quasar class (e.g. Bahcall et al. 1992) or to the Seyfert 1 class (e.g. Veron-Cetty & Veron 2001). Mrk 205 has a redshift z=0.07085 (Huchra et al. 1990) while NGC 4319 has z=0.00468 (Bowen et al. 1991).

In the standard paradigm of (radio-quiet) Active Galactic Nuclei (AGN), X-ray photons result from Compton upscattering of optical-UV photons in a hot thermal corona above the accretion disk surrounding a supermassive black hole. A large fraction of the seed photons are thought to be produced by thermal emission of the accretion disk itself, but Mrk 205 was found by McDowell et al. (1989) to have a weak blue bump. One of the possible scenarios which could explain this spectral morphology is that the bump may be highly variable and thus observed in a particularly weak state. In the thermal

comptonisation framework, blue bump variability would also induce hard X-ray variations.

A soft X-ray excess component may be interpreted as the hard tail of the blue bump (Walter & Fink 1993; Brunner et al. 1997). So far, evidence for a soft excess in Mrk 205 was only reported in two observations; with *EINSTEIN* IPC (Wilkes & Elvis 1987) and *XMM-Newton* (Reeves et al. 2001).

BeppoSAX, with its unique broad-band capabilities allows to extend the observation of Mrk 205 for the first time above 20 keV to study the properties of the hard X-ray emission as well as of a potential reflection component. Indeed, some of the X-ray photons can be reprocessed in the surroundings, in particular in the disk itself giving rise to neutral (or low-ionization) Fe K $\alpha$  line. While it is a common feature of Seyfert 1 galaxies (Nandra & Pounds 1994), only a few quasars present a clear detection of such a component (Williams et al. 1992; Nandra et al. 1997). Furthermore, the centroid of the line was often measured at energies close to 6.7 keV suggesting that its origin may be in an ionized layer of the accretion disk (Reeves & Turner 2000), the neutral Fe K $\alpha$  being emitted in distant and cold material lying outside the broad line region. Reeves et al. (2001), who analysed the XMM-Newton data of Mrk 205, claimed the existence of



**Fig. 1.** Summary of the X-ray observations of Mrk 205; evolution with time of the photon index Γ, the Hydrogen column density in the line of sight  $N_H$  (the dotted line denotes the Galactic value) and the 2-10 keV flux; units  $10^{-11}$  erg cm<sup>-2</sup> s<sup>-1</sup>. Data from Table 1. We have included as well an *IUE* light curve at 1300 Å; units  $10^{-14}$  erg cm<sup>-2</sup> s<sup>-1</sup> Å<sup>-1</sup>

a broad ionized component in addition to the neutral one in this object. However, Page et al. (2003) reprocessed the data including a reflection hump component in the fit and found little evidence for a broad, ionised component in the data.

In this paper, we present the analysis of three BeppoSAX observations of Mrk 205 for a total of 200 ks of exposure time. In Section 2, we summarize the knowledge we had on Mrk 205 prior to our BeppoSAX observation. In Section 3, we describe the data and their analysis while in Section 4 we compare our results and discuss them in the framework of the unified model.

#### 2. Previous X-ray observations of Mrk 205

Mrk 205 was observed by EINSTEIN HRI (Tananbaum et al. 1979: [a]; Zamorani et al. 1981), EINSTEIN IPC (Wilkes & Elvis 1987: [b]), EINSTEIN A-2 (Della Ceca et al. 1990), EXOSAT (Singh et al. 1991: [c]), GINGA (Williams et al. 1992: [d]), ASCA (Reeves & Turner 2000: [f]), ROSAT PSPC (Fiore et al. 1994:[e]; Fiore 1994; Rush et al. 1996b; Ciliegi & Maccacaro 1996; Fiore et al. 1997), ROSAT HRI (Arp 1996), XMM-Newton (Reeves et al. 2001: [g]; Page et al. 2003) and BeppoSAX (this work).

Mrk 205 was neither observed by RXTE nor Chandra to date.

We summarize the spectral parameters found in Fig. 1 and Table 1. The spectral variability was important over the 24 years spanned by the data (Fig. 1, upper panel). The photon indexes were measured mostly over the range  $2-10~{\rm keV}$  and thus were not affected by a soft excess component. However, the two ROSAT points (1991, 1992) are extrapolated values from the  $0.11-2.45~{\rm keV}$ , i.e. a spectral range very sensitive to the presence of a soft excess. They are certainly overestimated. In addition, several papers discuss the ROSAT PSPC calibration problems which results in a measure of the AGN X-ray spectral slopes being significantly steeper than those from other missions (Beckmann et al. 2002 and references therein).

We will discuss further in Section 4 the possible origin of the observed spectral variability in this object.

Little variations of the Hydrogen column density were observed during the same period (Fig. 1, second panel from the top). An excess of absorption may be present in the *XMM-Newton* 2000 observation, its reality will be discussed in Section 3.7.

The 2-10 keV flux seems to have decreased during the first  $\sim 18$  years of observation, a trend possibly correlated with a similar trend seen in the 1300 Å IUE light curve (Fig. 1, third panel from the top). Here again, the two ROSAT estimates of the 2-10 keV flux may be underestimated due to the overestimation of the 2-10 keV photon index. In the remainder of this section, we will focus on the XMM-Newton observation which took place one and a half year before our BeppoSAX pointings.

### 2.1. XMM-Newton observation

Reeves et al. (2001) reported an observation done with XMM-Newton on 6 May 2000. They measured a photon index  $\Gamma=1.86^{+0.02}_{-0.02}$  with an absorbing column in excess of the Galactic one limited to  $N_H<1.5\times10^{20}~\rm cm^{-2}$ . Fitting the MOS data only they found evidence for an ionized Fe line component at 6.7 keV along with the neutral Fe K $\alpha$  component. The same result was obtained using only the PN data, and thus the simultaneous fit of the data from both instruments gave tight constraints. They finally obtained the best fit with a narrow line resulting from neutral Fe K $\alpha$  emission at  $6.39\pm0.02~\rm keV$  ( $EW=56\pm23~\rm eV$ ) and a broad component at  $6.74\pm0.12~\rm keV$  ( $EW=135^{+65}_{-60}~\rm eV$ ). They did not detect Fe K edge, nor a reflection component but claimed that they could not rule out its presence

The observations were consistent with a relativistically-broadened accretion disk line model (Fabian et al. 1989) only if the disk ionization was high enough to produce Helike Fe at 6.7 keV and the neutral component was present. A clear soft excess was also detected, but fitting it with a black body gave a much too high temperature ( $kT=120\pm 8~{\rm eV}$ ) to be direct emission from the accretion disk. An explanation is enhanced emission from the upper layers of the ionized accretion disk. The fact that the

Table 1. Summary of the X-ray observations of Mrk 205; spectral parameters. The 2-10 keV flux was extrapolated from the best fit parameters

Mission, Reference	Date Γ		$N_H$	$Flux_{2-10keV}$	
			$10^{20} \text{ cm}^{-2}$	$10^{-11} \mathrm{erg} \ \mathrm{cm}^{-2} \ \mathrm{s}^{-1}$	
EINSTEIN (HRI), [a]	22-11-1978	$1.50^{f}$	$3.00^{+0.10}_{-0.10}$	$1.36^{+0.07}_{-0.07}$	
$EINSTEIN (IPC)^{\dagger}, [b]$	20-04-1980	$1.80^{+0.20}_{-0.20}$	$2.00^{+0.90}_{-0.80}$	$1.18^{+0.07}_{-0.05}$	
$EXOSAT^*$ , [c]	10-11-1983	$1.73^{+0.62}_{-0.55}$	$1.33^{+5.20}_{-1.30}$	$0.72^{+0.10}_{-0.11}$	
$EXOSAT^*$ , [c]	27-01-1984	$1.94^{+0.74}_{-0.41}$	$2.84^{+13.70}_{-2.35}$	$0.66^{+0.08}_{-0.09}$	
$GINGA^{\ddagger}, [d]$	16-01-1988	$1.95^{+0.13}_{-0.11}$	$3.40^{+3.40}_{-0}$	$1.01^{+0.10}_{-0.10}$	
ROSAT (PSPC), [e]	10-11-1991	$9.90 \pm 0.08$	$3.24^{+0.24}_{-0.24}$	$0.16^{+0.01}_{-0.01}$	
ROSAT (PSPC), [e]	15 - 04 - 1992	$2.33_{-0.08}^{+0.10}$ $2.32_{-0.10}^{+0.10}$	$3.32^{+0.30}_{-0.30}$	$0.11^{+0.01}_{-0.01}$	
$ASCA^{\ddagger}, [f]$	03-12-1994	$2.13^{+0.03}_{-0.03}$	$2.50^{+0.90}_{-0.90}$	$0.75^{+0.09}_{-0.09}$	
BeppoSAX, this work	12 - 03 - 1997	$1.77^{+0.07}_{-0.06}$	$2.90^{+1.18}_{-0}$	$1.21^{+0.02}_{-0.02}$	
$XMM$ - $Newton^{\dagger \ddagger}$ , [g]	06-05-2000	$1.86^{+0.02}_{-0.02}$	4.30	$0.47^{+0.05}_{-0.05}$	
$BeppoSAX^{\dagger \ddagger}$ , this work	21-12-2001	$1.87^{+0.02}$	$2.90^{+0.25}_{-0}$	$1.11^{+0.01}$	
BeppoSAX, this work	05-01-2002	$1.72^{+0.07}_{-0.07}$	$3.69^{+2.75}_{-0.79}$	$0.72^{+0.01}_{-0.01}$	

Note: Galactic  $N_H = 2.90 \times 10^{20} \text{ cm}^{-2}$  (Dickey & Lockman 1990)

Compton reflection from ionized material model (PEXRIV: Magdziarz & Zdziarski 1995) gave good results (with  $\xi = 300\pm80~{\rm erg~cm~s^{-1}}$ ) was probably confirming this (see also Nayakshin et al. 2000). No warm absorber was found; the limits on the OVII and OVIII edges were  $\tau < 0.12$  and  $\tau < 0.15$  respectively.

Page et al. (2003) reprocessed these data and found that an acceptable fit was obtained with a power law plus neutral emission line at 6.4 keV. They noticed that a better fit was obtained either with a Compton reflection component or with an ionized emission line, with a similar goodness of fit. The presence of ionized relativistic material was thus not unambiguously detected in this object, the illumination of a distant cold material providing a simpler self-consistent explanation.

#### 3. BeppoSAX observation

Mrk 205 was observed by BeppoSAX (Boella et al. 1997a) on 21 December 2001 (LECS: 47582 s, MECS: 165458 s, PDS: 70549 s) and on 5 January 2002 (LECS: 3611 s, MECS: 32603 s, PDS: 12606 s). In addition, public data from the archive (12 March 1997, LECS: 2525 s, MECS: 14727 s, PDS: 6646 s; part of the core program and pointed on NGC 4319) are included in the present study.

LECS (Low Energy Concentrator Spectrometer, Parmar et al. 1997) and MECS (Medium Energy Concentrator Spectrometer, Boella et al. 1997b) data were reduced following the standard procedures. For PDS (Phoswich Detector System, Frontera et al. 1997) data, high level products from the Sax Data Center were used. LECS and MECS data have been screened according to Fiore et al. (1999) to produce equalized event files using an extraction

radius of 8' for LECS and 4' for MECS. Data from the two MECS units have been merged to increase the signal-to-noise ratio.

For the longest exposure (2001), the average source flux in the 2-10 keV band (MECS data only, power law model with Galactic absorption) was  $(1.11\pm0.01)\times10^{-11}$  erg cm<sup>-2</sup> s<sup>-1</sup> corresponding to a luminosity  $L=2.55\times10^{44}$  erg s<sup>-1</sup> in the 2-10 keV band rest frame (here and in the following  $H_0=50$  km s<sup>-1</sup> Mpc<sup>-1</sup>). From that luminosity, a black hole mass of about  $2\times10^6~M_{\odot}L_{Edd}/L$  could be inferred. These values are more typical of a luminous Seyfert 1 galaxy than a quasar. The 2-10 keV flux measured in 2001 was twice the value measured one and a half year before by XMM-Newton (Fig. 1).

# 3.1. Image analysis: unprobable influence from NGC 4319 and other sources

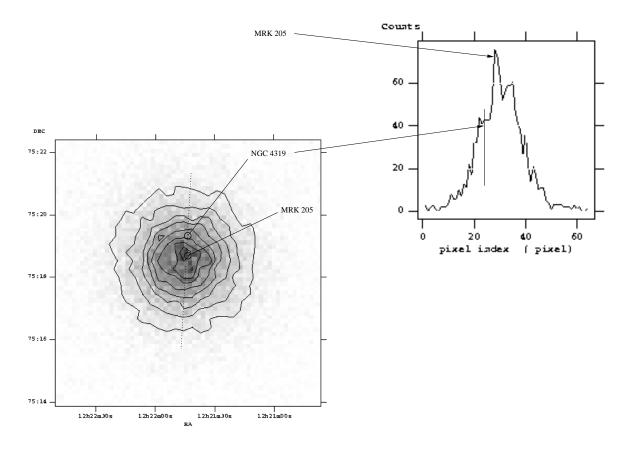
We analysed the MECS  $1.3-10~\rm keV$  image (angular resolution 42'') of the 2001 observation and compared it with a  $ROSAT~\rm HRI~0.1-2.4~\rm keV$  image (angular resolution of 5''). For the HRI image, Arp (1996) evaluated the shift of the position from the optical to X-ray to be of at most 7'', equivalent to one MECS pixel. Three sources were identified in the MECS image, Mrk 205 included. Being in the field of view of the PDS also (1.4 degrees while the MECS field of view is 28'), we investigated if one of the other two sources could contribute to the Mrk 205 PDS spectra. In addition, one had to make sure that NGC 4319 located at 40'' from Mrk 205 (and thus not resolved by this instrument) was not emitting in the X-rays (we used the following J2000 coordinates for Mrk 205:  $12h~21'~44.12''~+75^\circ~18'~38.2''$  and for NGC 4319:

 $<sup>^{\</sup>dagger}$  evidence for a soft excess component reported

<sup>\*</sup> broken power law model preferred

 $<sup>^{\</sup>ddagger}$  evidence for an Fe K $\alpha$  line reported

f fixed parameter



**Fig. 2.** Left panel: MECS intensity image of the 2001 observation with contours at (from the outside) 7.7, 15.5, 23.3, 31.1, 38.8, 46.6, 54.4, 62.2 counts. Right panel: profile along the dotted line displayed on the left panel (1 MECS pixel is 7"). NGC 4319 is located 6 pixels on the left of the central peak in a region which accounts for 54% of the peak value

12h 21' 44.07" +75° 19' 21.3", these are located by circles on Fig. 2). We discuss those points below.

- According to the profile displayed in Fig. 2 (right panel), NGC 4319 is located in a zone in which half of the MECS peak count rate is received. Furthermore the shape itself does not allow to tell if NGC 4319 is really detected (the break in the profile can be a feature of the MECS PSF). It was not detected by ROSAT HRI (Arp 1996) in the soft X-ray domain, but it could be Compton thick with emission only detected above 5-10 keV. However, there are no radio sources (apart from Mrk 205) above 2.5 mJy within a radius of 1.5' of Mrk 205 in the NVSS catalog (Condon et al. 1998). Thus it is very unlikely that NGC 4319 is an AGN;
- 2. We have fitted the 2001 MECS spectrum with a power law, extrapolated the model in the PDS domain and added the PDS data without fitting them but including an intercalibration constant MECS/PDS frozen to the recommended value of 0.88 (see Fiore et al. 1999).

- This is of course not taking into account any reflection component. Fig. 3 shows the resulting spectra along with the residuals of both data sets with respect to the MECS best fit model. It is very unlikely that emission from NGC 4319 contributes to the excess at 20 keV and matches the extrapolation of the MECS flux at high energies.
- The other MECS 3. sources in the are  $J122206.4+752613^{1}$ XMMU which is a narrow emission line galaxy and NGC 4291, a non-active elliptical galaxy. The first has a MECS count rate 10 times smaller than Mrk 205 while for the latter it is 100 times smaller. According to Fiore et al. (1999), those sources are certainly in the residual PDS count rate (estimated by comparing to the known MECS count rate) obtained in blank fields (Guainazzi & Matteuzzi 1997);

 $<sup>^{1}</sup>$  the source XMMU J122206.4+752613 corresponds to MS 1219.9+7542 and to ROSAT 1RXS J122214.2+752624

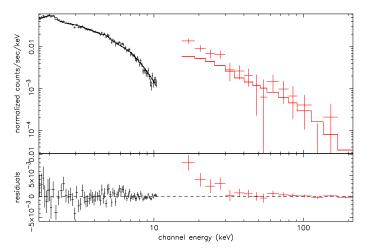


Fig. 3. Upper panel: fit of the 2001 MECS data with a power law and absorption frozen to the Galactic value. PDS data are included without being fitted but with an intercalibration constant frozen to 0.88. Lower panel: residuals with respect to the MECS best fit model

From these different arguments, we conclude that neither NGC 4319 nor one of the two other sources detected in MECS contribute significantly to the X-ray spectrum of Mrk 205.

# 3.2. Possible absorption in the spiral arms of NGC 4319

Of type  $S_{pec}$ , NGC 4319 is an extremely disrupted spiral galaxy (Arp 1996) with an ISM very different from the Milky Way's ISM. The ISM of NGC 4319 has been thoroughly studied by various authors which have found Ca II K and Mg I b absorption features at the redshift of NGC 4319 (Stockton et al. 1979).

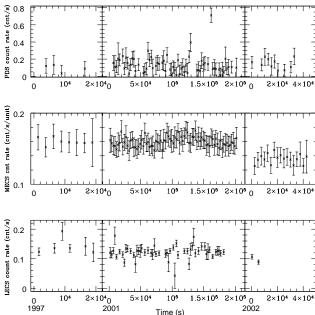
Bahcall et al. (1992) succeeded in the detection of Mg II doublet in NGC 4319 and deduced an approximative H I column density of about  $3\times 10^{18}~\rm cm^{-2}$  from the Mg II column density, precised by Womble et al. (1993) to be N(H I)  $< 5.2\times 10^{19}~\rm cm^{-2}$ .

The gas in the inter-arm region of NGC 4319 shows a higher grade of ionisation than in the Milky Way and the unusually low amount of HI (or Mg II) is not understood for this kind of late-type galaxy.

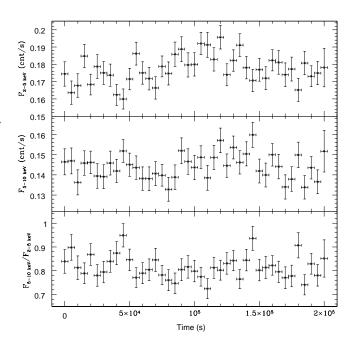
The X-ray absorption in the  $0.5-2.0~\rm keV$  energy range is however dominated by CNO K-shell absorption, but no evidence for an enhanced abundance of heavy elements is reported in the literature. To conclude, we do not find evidence for absorption in the Mrk 205 data due to the presence of NGC 4319 in the line of sight.

#### 3.3. Variability

We produced light-curves from the event files using Xronos 5.16 in a single binning of 2500 s to ease intercomparison between the different observations (Fig. 4). Note that 15 days separate the middle panel of Fig. 4 to the right panel. LECS and MECS light curves of the



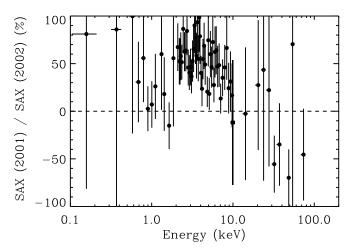
**Fig. 4.** Light curves of the 3 observations. Top panels: PDS 13-200 keV, middle panels: MECS 2-10 keV (count rate per MECS unit), bottom panels: LECS 0.01-2 keV



**Fig. 5.** Observation of December 2001, MECS data. Top panel: 2-5 keV light curve, middle panel: 5-10 keV light curve, bottom panel: hardness ratio

statistically best observation (2001) were surprisingly different. Cross-correlation analysis gave a weak correlation coefficient between these two light curves.

From 1997 to 2001, the LECS and MECS count rates remained constant. Since a MECS unit was lost between those 2 observations, we have plotted in Fig. 4 the MECS



**Fig. 6.** Spectral variations between 2001 and 2002 (ratio of the 2001 spectra to the 2002 spectra)

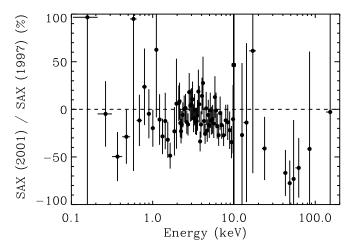


Fig. 7. Spectral variations between 1997 and 2001 (ratio of the 2001 spectra to the 1997 spectra)

light curves for a single unit. The average PDS count rate seems to have increased by about 40%, but this latter value is very inaccurate due to the few points in 1997. From 2001 to 2002, the LECS average count rate seemed to have dropped by about 16% percent, but this latter value is based on a small fraction of the observation time. The MECS average count rate dropped by 15%, while the PDS count rate remained constant.

From the December 2001 observation (Fig. 4, middle panels; see also Fig. 5) one sees that no significant variability was observed on very short timescales (of the order of hours). On the other hand, long term variability (on week and year timescale) can be quite important (Fig. 1).

In Fig. 5, we display the light curves  $2-5~{\rm keV}$  and  $5-10~{\rm keV}$  along with the hardness ratio, for the longest observation (2001). The  $2-5~{\rm keV}$  energy band showed a weak minimum (10% below the average flux) around  $4\times10^4~{\rm s}$  which was not present in the  $5-10~{\rm keV}$  range. The result-

ing hardness ratio thus showed an hardening around that time. The differences in the two light curves are certainly not due to  $N_H$  variations as a value of the order of  $10^{20}$  cm<sup>-2</sup> is unlikely to change the spectra above 2 keV.

# 3.4. Model independent spectral analysis

To study the spectral variations between the different observations, we have divided the 3 spectra (LECS, MECS, PDS) of the 2001 observation from the 3 spectra of the 2002 observation (Fig. 6). We have also divided the 3 spectra of the 2001 observation from the 3 spectra of the 1997 observation (Fig. 7). We have taken the ratios for each energy bins. One sees an important variation in the  $2-10~{\rm keV}$  energy range between the 2001 and 2002 observation which did not occur between the 1997 and 2001 observation, in agreement with Fig. 4. A rough comparison of Fig. 6 and Fig. 7 shows that the LECS and PDS points may have been enhanced as well in Fig. 6, but the signal-to-noise is not good enough to conclude that the soft  $0.1-2.0~{\rm keV}$  or hard  $10-220~{\rm keV}$  bands varied significantly.

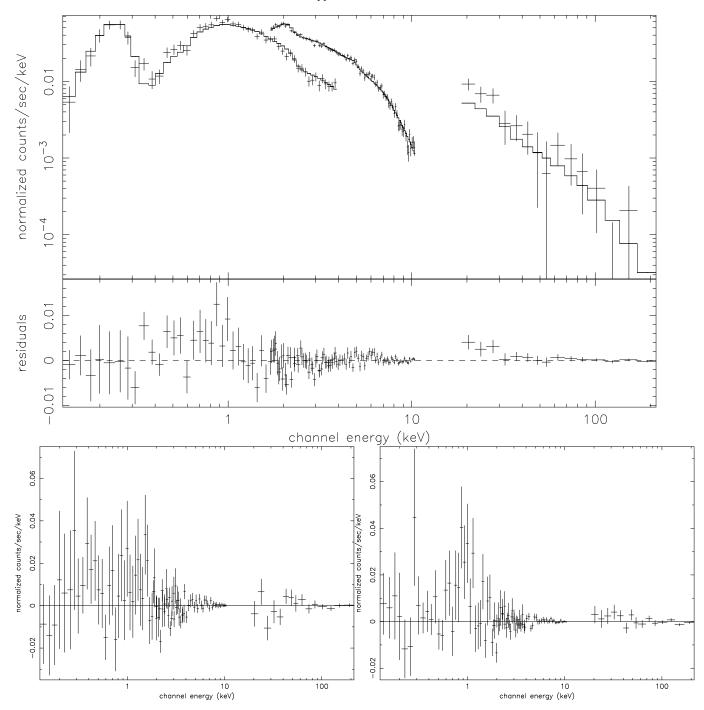
# 3.5. Spectral analysis

The spectral fitting was performed with XSPEC 11.2.0. In the following, all errors are quoted at the 90% confidence level for one interesting parameter. We have used the recommended energy bounds (LECS: 0.12-4.0 keV, MECS: 1.65-10.5 keV, PDS: 15-220 keV), and the intercalibration "constants" were free to vary in the intervals 0.7-1.0 and 0.82-0.90 for LECS/MECS, PDS/MECS respectively (Fiore et al. 1999).

The upper panel of Fig. 8 represents a fit of the 2001 data with a power law model modified by interstellar photoelectric absorption. The photoelectric absorption cross sections used were those of Morrison and McCammon (1983).

The lower panels of Fig. 8 show the fits residuals with the same model for the 1997 and 2002 observations respectively. The best fit spectral parameters found are presented in Table 2. The fit of the 2001 observation is poor with an excess in the  $15-40~{\rm keV}$  range as well as in the soft X-rays, up to about 1 keV, finally a small feature around 6 keV is present.

For the 1997 and 2002 observations, the fits are already acceptable with this simple model for the continuum. In the 2002 observation, a line-like feature is clearly detected around 1 keV while a small excess in the range 15-40 keV may be present. The signal-to-noise ratio of the continuum is not good enough to accurately constrain more complex physical models.



**Fig. 8.** Upper panel: broad band (folded) spectrum of the 2001 observation. The spectrum is fitted with a power law plus Galactic absorption model. Lower panel, left side: residuals after the fit of the 1997 observation with the same model. Lower panel, right side: as above, for 2002

#### 3.6. Spectral analysis: reflection component

We started a more detailed analysis by fitting the 2001 observation data above  $\sim 2$  keV, i.e. using MECS and PDS data only. As  $N_H$  was not constrained in this case, we fixed it to the Galactic value. We tried 2 different models for the continuum shape.

The first was a cut-off power law plus a Compton reflection from dense, neutral matter (PEXRAV: Magdziarz & Zdziarski 1995). This model is meant to describe the

effects of Compton scattering of photons from an X-ray source associated with the inner parts of the accretion disk. The spectrum of reflected X-ray was computed as a function of disk inclination using the transfer function of Magdziarz & Zdziarski (1995). The free parameters of this model are the spectral index  $\Gamma$  of the primary power law, the folding (upper cutoff) energy  $E_c$  of the primary X-ray spectrum and the normalization of the reflection component R. In our fits, the inclination angle of the disk was frozen to 30 degrees as it is usually admitted for

Table 2. Fits of the BeppoSAX data using a power law plus Galactic absorption model: best fit spectral parameters

Observation	Γ	$N_H$	$\chi^2(dof)$ , proba		
		$10^{20} \text{ cm}^{-2}$			
2001	$1.87_{-0.03}^{+0.02}  1.77_{-0.06}^{+0.07}  1.72_{-0.07}^{+0.07}$	$< 3.15^p$	182.23(139), 0.008		
1997	$1.77^{+0.07}_{-0.06}$	$< 4.08^{p}$	106.88(99), 0.277		
2002	$1.72^{+0.07}_{-0.07}$	$3.69^{+2.75}_{-0.79}$	104.70(99), 0.328		

<sup>&</sup>lt;sup>p</sup> pegged at Galactic value  $N_H = 2.90 \times 10^{20} \text{ cm}^{-2}$ 

Table 3. Reflection component: fits of the MECS and PDS data (1.65-220 keV). A stands for photoelectric absorption, PL for power law, PEXRAV for the Compton reflection from cold material and COMPHA for the thermal comptonisation model.  $\Gamma$  is the photon index of the primary power law,  $E_c$  is the folding (upper cutoff) energy, R is the normalization of the reflection component,  $kT_e/m_ec^2$  is the corona temperature (normalized to the electron mass),  $\tau$  is the corona optical depth,  $E_{\rm Fe}$  is the Fe K $\alpha$  line energy (rest frame) and  $EW_{\rm Fe}$  its equivalent width. As  $N_H$  is not constrained above 1.65 keV, we fixed it to the Galactic value. In addition, the results of a fit of the XMM-Newton data above 2 keV are presented

Obs.	Model	Γ	$E_c$	R	$kT_e/m_ec^2$	τ	$E_{\mathrm{Fe}}$	$EW_{\mathrm{Fe}}$	$\chi^2(dof)$ , prob.
\ <u></u>			keV				keV	$\mathrm{eV}$	_
SAX 2001	A+PL	$1.83^{+0.03}_{-0.03}$	-	-	-	-	-	-	137.47(93), 0.002
	A+PEXRAV+G	$2.00^{+0.03}_{-0.03}$	$487.17^{+u}_{-304.40}$	$1.74^{+0.25}_{-0.32}$	-	-	$6.51^{+0.42}_{-0.55}$	71.30	100.38(89), 0.192
	A+COMPHA+G	-	-	$3.22^{+0.99}_{-1.35}$	$0.32^{+0.04}_{-u}$	$0.15^{+0.33}_{-0.08}$	$6.52^{+0.18}_{-0.18}$	70.30	100.39(89), 0.192
SAX 1997	A+PL	$1.74^{+0.09}_{-0.08}$	=	-	-	-	-	-	48.36(53), 0.655
	A+PEXRAV+G	$1.77^{+0.70}_{-0.13}$	$240.00^{f}$	< 8.78	-	-	$6.40^{f}$	<1	48.16(51), 0.587
	A+COMPHA+G	-	-	$2.50^{f}$	$0.48^{+0.04}_{-u}$	$0.05^{+u}_{-u}$	$6.40^{f}$	<1	47.66(51), 0.607
SAX 2002	A+PL	$1.69^{+0.07}_{-0.07}$	-	-	-	-	-	-	49.91(53), 0.595
	A+PEXRAV+G	$1.80^{+0.25}_{-0.16}$	$228.30^{+u}_{-193.00}$	< 3.47	-	-	$6.40^{f}$	34.90	46.89(50), 0.599
	A+COMPHA+G	-	-	$2.43^{+4.41}_{-1.97}$	$0.37^{+0.13}_{-0.09}$	$0.19^{+u}_{-u}$	$6.40^{f}$	38.90	46.27(50), 0.624
XMM 2000	A+PEXRAV+G	$1.73^{+0.14}_{-0.14}$	$500.00^{f}$	$0.48^{+1.09}_{-0.48}$	-	-	$6.40^{f}$	75	75.24(54), 0.029

 $<sup>^</sup>f$  parameters fixed in the fit

Seyfert 1 galaxies, while the abundances of Fe and other heavy elements were fixed to the solar values.

The second model was a more physical model of thermal Comptonisation for a disk-corona configuration in a slab geometry (COMPHA: Haardt 1993; Haardt et al. 1997). The table model for XSPEC<sup>2</sup> includes both the primary continuum and the reflection. The reflection component was computed following White et al. (1988) and Lightman & White (1988) assuming reflection on neutral matter and a constant spectral shape (angle-averaged) for the reflected photons. The free parameters of the model are the corona temperature normalized to the electron mass  $kT_e/m_ec^2$ , the corona optical depth  $\tau$  and the normalization of the reflection component R (without the comptonisation effects). In our fits, the temperature  $kT_{bb}$  of the disk (which is assumed to radiate like a single black-body) was frozen to 10 eV while the inclination angle of the disk was frozen to 30 degrees.

Both continuum models were completed by the addition of a Gaussian emission line to model the Fe K $\alpha$  line with a width of 0.1 keV frozen (a more detailed analysis on the Fe line parameters is made below). The best fit spectral

The  $\chi^2$  is not significantly improved by PEXRAV for the two other observations (Table 3). We thus cannot constrain the presence of this component in 1997 and 2002.

The fit of the December 2001 observation is improved by using COMPHA when compared to a power law ( $\Delta\chi^2{=}30.11$  for 2 more parameters; F-test value 12.80 and probability  $1.30\times10^{-5}$ ). It is slightly improved by including the Fe line ( $\Delta\chi^2{=}7.17$  for 2 more parameters; F-test value 3.19 and probability  $4.61\times10^{-2}$ ).

To summarize, fitting with the thermal comptonisation model COMPHA does not improve the fit of the 2001

<sup>&</sup>lt;sup>u</sup> unbounded parameters

parameters are given in Table 3. The fit of the observation made in 2001 is improved by using PEXRAV without cutoff  $(\Delta\chi^2{=}29.51$  for 1 more parameters; F-test value 25.20 and probability  $2.57\times 10^{-6}).$  It is slightly improved when including the cutoff  $(\Delta\chi^2{=}1.97$  for 1 more parameters; F-test value 1.70 and probability 0.196) and the line  $(\Delta\chi^2{=}5.60$  for 2 more parameters; F-test value 2.48 and probability  $8.93\times 10^{-2}).$  The solid angle subtended by the reflector is poorly constrained (Fig. 9). Figure 10 shows that the cutoff upper limit is not constrained.

<sup>&</sup>lt;sup>2</sup> http://pitto.mib.infn.it/~haardt/ATABLES

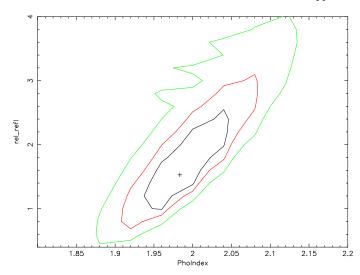


Fig. 9. Contour plot R- $\Gamma$  after the fit of the 2001 observation with the Compton reflection from cold material model plus an Fe K $\alpha$  line

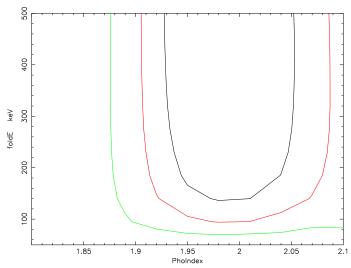


Fig. 10. Contour plot  $E_c$ - $\Gamma$  after the fit of the 2001 observation with the Compton reflection from cold material model plus an Fe K $\alpha$  line

observation in comparison to the Compton reflection from cold material PEXRAV fits.

## 3.7. Spectral analysis: soft excess

We now consider the broad band 2001 spectrum, adding the LECS 0.12-4.0 keV data to the MECS and PDS data. Freezing all the parameters of the best fit PEXRAV model above 2 keV and fitting again the broad band spectrum, we get an  $N_H = (3.46^{+0.26}_{-0.22}) \times 10^{20}$  cm<sup>-2</sup>, and a  $\chi^2$  of 138.49 for 142 degrees of freedom (dof). We found no clear evidence for a soft excess of photons in the LECS, even if the intercalibration constant is set to its lowest value of 0.7. To have an upper limit on a soft excess component, we add a black body to our model, fixing its temperature to the one obtained by Reeves et al. (2001) with the XMM-Newton data, i.e. kT = 120 eV. We found an

 $N_H = (3.60^{+0.36}_{-0.29}) \times 10^{20}$  cm<sup>-2</sup> and a  $\Delta \chi^2 = 1.73$  for 1 more parameter; F-test value 1.78 and probability 0.184. The soft excess is here much weaker ( $\sim 3\%$  of the emission over the range 0.5-2.0 keV) than it was claimed in the *XMM-Newton* observation ( $\sim 15\%$  of the emission in the range 0.5-2.0 keV).

We re-analysed the XMM-Newton PN data to see whether the soft excess component was as important as claimed by Reeves et al. (2001) when the reflection component was present. Indeed, Reeves et al. (2001) discussed the soft excess component modeling the continuum with a power law while Page et al. (2003) used the Compton reflection model but stopped their analysis at 3 keV. As before, we deduced the parameters of the continuum fitting the data above 2 keV (Table 3). The PEXRAV model produced a  $\chi^2$ of 75.24 (54 dof) with a realistic reflection parameter R of 0.48. We then fixed the continuum and included the 0.5 to 2.0 keV points. The residuals showed a clear excess below 2 keV. Adding a black body with a temperature of 120 eV did not produce a good fit ( $\chi^2 = 463.29$  for 232 dof) but, when this parameter was released, the fit was much better and the temperature significantly increased (see Table 4 and Fig. 11). We did not find an unusual value of  $N_H$ (Fig. 1) anymore.

The contribution of the soft excess component is thus more important than previously claimed: it represents  $\sim 23\%$  of the emission in the range 0.5 to 2.0 keV.

However, the black body parameters do show suspicious values meaning that the emission is probably not a simple black body.

## 3.8. Spectral analysis: Fe K $\alpha$ line

To assess the importance of an Fe K $\alpha$  line in our data, we concentrate on the MECS data. In this case,  $N_H$  is not constrained, we thus fixed it to the Galactic value. Using only the power law model, it is clear from the residuals in the lower panel of Fig. 8 that a line feature is detected in the 2001 observation, while it does not seem to be present in the 2002 observation (same figure). The statistics of the 1997 observation is too poor to give conclusive results.

On the other hand, after fitting the continuum of the 2001 data with the Compton reflection model, the residuals do not show a clear evidence for a line feature anymore.

We thus concentrate on the 2001 best fit observation and add a Gaussian emission line. The fit is improved ( $\Delta\chi^2=6.07$  for 1 more parameter; F-test 5.39 and probability  $2.29\times 10^{-2}$ ) with a line at 6.53 keV consistent with neutral Fe K $\alpha$  and weakly ionized species. Note that the line is not resolved as  $\sigma_{Fe}<500$  eV is smaller than the MECS energy resolution.

Both Compton reflection from cold material and thermal comptonisation model give a line energy of 6.53 keV but the latter gives a slightly larger equivalent width. According to George & Fabian (1991), if a reflection is

**Table 4.** Soft excess: fits of the LECS, MECS, PDS data and comparison with XMM-Newton data. A stands for photoelectric absorption, PEXRAV for the Compton reflection from cold material, G for a Gaussian line and BB for a black body. We first consider the spectra above 2 keV, freezing all the parameters of the continuum  $(\Gamma, R, E_c)$  to their best fit value of Table 3. We then add the data below 2 keV, add a soft excess component, and fit again the broad band spectrum

Obs.	Model	$N_H$	$kT_{\mathrm{BB}}$	$\chi^2(\mathrm{dof})$	proba
		$10^{20} \text{ cm}^{-2}$	${ m eV}$		_
BeppoSAX 2001	A+PEXRAV+G+BB	$3.60^{+0.50}_{-0.31}$	$100.00^{+u}_{-99.90}$	137.40(140)	0.546
XMM 2000	A+PEXRAV+G+BB	$< 3.14^{p}$	$164.86^{+4.85}_{-8.57}$	325.09(231)	< 0.001

 $<sup>^{</sup>u}$  unbounded parameters

<sup>&</sup>lt;sup>p</sup> pegged at Galactic value  $N_H = 2.90 \times 10^{20} \text{ cm}^{-2}$ 

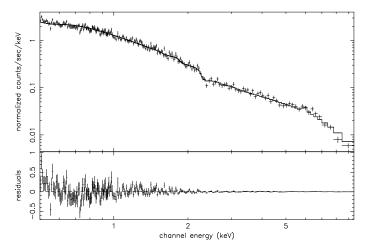


Fig. 11. XMM-Newton 2000 data, fit with the Compton reflection from cold material model

**Table 5.** Fe K $\alpha$  line: fits of the 2001 observation, MECS data only. A stands for photoelectric absorption (fixed to the Galactic value), PEXRAV for the Compton reflection from cold material, COMPHA for the thermal comptonisation model and G for Gaussian (line energies are rest frame). R,  $E_c$ ,  $kT_e/m_ec^2$  and  $\tau$  were fixed to their best fit values of Table 3

Model	Γ	$E_{\mathrm{Fe}}$	$\sigma_{ m Fe}$	$EW_{\mathrm{Fe}}$	$\chi^2(\mathrm{dof})$	prob
		keV	${ m eV}$	$\mathrm{eV}$		
A+PEXRAV	$1.99^{+0.03}_{-0.02}$	-	-	-	92.81(78)	0.121
A+PEXRAV+G	$1.99^{f}$	$6.53^{+0.23}_{-0.19}$	< 447.60	56.80	86.74(77)	0.210
A+COMPHA	-	-	-	-	95.53(79)	0.099
A+COMPHA+G	-	$6.53^{+0.16}_{-0.16}$	< 314.82	71.90	86.25(77)	0.220

f parameters fixed in the fit

present, it has to be associated with an Fe line. In the neutral case, they computed an equivalent width of the order of 1.7 keV with respect to the reflection or  $\sim 150$  eV with respect to the continuum, for R=1 (see also Zdziarski et al. 2002). Our measures of the line equivalent width and R are inconsistent with the predictions of George & Fabian (1991). We shall discuss in Section 4 a possible reason to this deviation.

Unlike Reeves et al. (2001) and Page et al. (2003), no residuals are found at  $\sim 7$  keV, but the line equivalent width and centroid suggest that it could also be a blend of a neutral narrow component and a weakly ionized one. Adding a second Gaussian to the model while freezing the

line energies to 6.4 keV and 6.7 keV rest frame following Reeves et al. (2001) was not significant and neither was the relativistically-broadened accretion disk line model (Fabian et al. 1989) which was preferred by Reeves et al. (2001). For the 2002 observation, the residuals were not improved when using the Compton reflection model and thus were identical to Fig. 8.

#### 3.9. Line feature at 1 keV

The spectrum of the 2002 observation contains a narrow emission line feature centered at about 1 keV (see Fig. 8, lower right panel). We have fitted a model composed of a power law plus Galactic absorption and a Gaussian com-

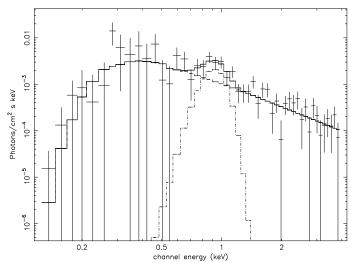


Fig. 12. Observation taken in 2002, LECS data only, unfolded spectrum fitted with a power law plus Galactic absorption and a Gaussian component centered at an energy of about 1 keV

ponent (Fig. 12) and found the following results. The line is centered at  $0.98^{+0.08}_{-0.34}$  keV (quasar rest frame) has a width of  $0.11^{+0.33}_{-0.11}$  keV (not resolved; the LECS energy resolution at 1 keV is  $\sim 0.2$  keV) and an equivalent width of 518 eV. The addition of the line significantly improves the fit at better than the 98% confidence level according to the F-test. Note however that even if this simple model allows us to estimate the line centroid energy and total flux, it is physically unrealistic as it is likely that this line results from a blend of several emission lines, predominantly from ionized species of Ne and possibly from Fe L shell.

A similar feature was also found in the ASCA spectrum of the quasar PG 1244+026 (Fiore et al. 1998; George et al. 2000) as well as in the ASCA spectrum of the Seyfert 1 galaxies TON S180 (Turner et al. 1998) and Akn 564 (Turner et al. 1999).

We searched the *XMM-Newton* data for this feature but did not find a clear evidence for it (the equivalent width reached a maximum of 10 eV).

# 4. Discussion-Conclusion

The nature of Mrk 205. It has been noted that, in contrast to the Seyfert 1 case, the detection of Fe K $\alpha$  line in quasars was more seldom and that if detected, the line energy was around 6.7 keV instead of 6.4 keV (Reeves & Turner 2000). It was thus suggested that sensitive instruments like the EPIC camera on XMM-Newton could observe the Fe K $\alpha$  emission of quasar in detail. From our X-ray observations of Mrk 205, and from the luminosity we obtained from them, we suggest to be careful in classifying this object. Its luminosity is found in the region in which it is difficult to choose for a Seyfert 1 or for a quasar. It seems thus difficult to generalize any findings made on this object to a whole class.

The soft excess variability. The soft excess seems to be variable in this source although a very strong soft excess was never reported. It seems to be absent/not detected in 1997, strong in 2000, weak in 2001 and absent again in 2002 (thus varied in a 2 weeks timescale). The variability of the soft excess cannot be associated with variability of a cold absorber as we have shown in Fig. 1 and Section 3 that the absorbing column density deduced from spectral modeling was roughly constant over all the observations.

Note that the soft excess component measured here should not be associated with direct thermal emission from the accretion disk because its temperature is much too high.

Thermal comptonisation. The historic variations of the photon index can be explained in the framework of thermal comptonisation. Mrk 205 showed large 2-10 keV flux variations between the 1997 (high flux state), 2000 (low flux state), 2001 (high flux state) and 2002 (low flux state) observations which do not seem to be correlated with variability of the soft excess.

However, in thermal comptonisation processes, the index of the comptonised component is inversely correlated with the ratio of the photon seed flux to the comptonised power law flux. We can test this hypothesis if we assume that the soft excess strength is proportional to the strength of the soft photon field (Walter & Fink 1993; Walter & Courvoisier 1990). The index for 2000 was  $1.73 \pm 0.02$  for a ratio of 23% while for 2001 the index was  $2.00 \pm 0.03$  for a ratio of 3%. This is thus in disagreement with our expectations.

On the other hand, using the COMPHA model to estimate the UV flux, we found that an increase of the photon index goes with a larger UV flux, in agreement with simple thermal comptonization model expectation. Similar conclusions were found by Petrucci et al. (2000) for the Seyfert 1 galaxy NGC 5548. The contradiction between the soft X-ray and UV variations is not very surprising. Indeed the soft excess component is certainly linked to the UV bump through complex radiative transfer (e.g. comptonization in the warm (few keV) upper layer of the accretion disk). Consequently, its expected spectral behavior is not straightforward and may be at odds from the zeroth order expectations.

Constraints on the reflecting material. The nature of the reprocessing medium can be the accretion disk or some obscuring, distant material. If the Fe K $\alpha$  line is produced in the inner parts of the accretion disk, its profile should be considerably broader, with very distinct asymmetries due to Doppler and gravitational redshift. This is not possible to test with the BeppoSAX data while Page et al. (2003) report that they did not find evidence for it in the XMM-Newton data either.

The strength of the Compton reflection R is here systematically higher than what was found in Seyfert galaxies  $(R \simeq 0.8 \pm 0.2, \text{ Nandra \& Pounds 1994})$ , while on the

contrary, the iron line equivalent width EW is smaller than the mean EW value. It is worth noting however that our measures of R are consistent with R < 1 at 90% confidence, see Fig. 9.

The inconsistency between the best fit values of R and the line EW may then occur by chance and the reflecting material in Mrk 205 may simply cover a solid angle  $< 2\pi$  as seen from the X-ray source. The line energy then also suggests the material to be neutral. A low Fe abundance will make the line weaker and the Compton reflection hump stronger (as observed) by reducing the opacity above the Fe K edge (George & Fabian 1991).

Given these constraints a simple solution could be that the reprocessing is done in a distant material e.g. the dust torus.

However, we cannot rule out the presence of close mildly ionized material. Indeed Nayakshin & Kallman (2001) have shown that in the case of relatively strong illumination by local X-ray flares, the upper layer of the reflecting material may become strongly ionized. In this case the reflection features are mainly produced by the neutral depth layer and the expected line energy is 6.4 keV as we observed. The line EW is however smaller than the one expected in the case of neutral reflector since the line may be partly comptonized in the hot skin. But clearly this skin cannot be highly ionized since in this case it would act like a perfect mirror and the reflection hump would be suppressed contrary to what we observe. Lines in the soft band are also expected in the case of a mildly ionized medium and it may explain the poor fit of the soft excess with a simple black body component.

The 1 keV line. The spectrum of January 2002 contains a narrow emission feature centered at an energy of  $\sim 1~\rm keV$  which was also found in at least 3 other AGN. This feature could be explained by a blend of several emission lines, predominantly from ionized species of Ne and possibly from Fe L shell. The line emission could be a result of reflection in photoionized matter. It could also be resulting from recombination in an optically thin thermal plasma.

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